

Imperfectly grown periodic medium: absence of localized states

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Abstract

We consider a discrete model of the d -dimensional medium with Hamiltonian $\Delta + v$; the lattice potential v is constructed recursively on a nested sequence of cubes Q_n obtained by successive inflations with integer coefficients. Initially, the potential is defined on the cube Q_0 . At the n th step the potential, which is already constructed on the cube Q_{n-1} , gets extended Q_{n-1} -periodically to the cube Q_n ; then its values at m_n randomly chosen points of Q_n are arbitrarily changed. This alternating process of periodic extension and introduction of impurities goes on, resulting in an (in general, unbounded) potential v . We show that if the size of the cube Q_n grows fast enough with n while the sequence m_n grows not too fast, then the Schrödinger operator $\Delta + v$ almost surely does not have eigenvalues.

1 Introduction

We consider a discrete model of the d -dimensional medium with Hamiltonian $h_v = \Delta + v$. Here Δ is the lattice Laplacian, and the lattice potential $v(\cdot)$ is the result of the following recursive process. Consider a nested sequence of cubes Q_n ($n = 0, 1, 2, \dots$) of the form

$$Q_n = \{x \in \mathbf{Z}^d: -\frac{T_n}{2} < x_i \leq \frac{T_n}{2}, \quad i = 1, \dots, d\} \quad (1)$$

whose sides $T_n \in \mathbf{N}$ are such that

$$\frac{T_n}{T_{n-1}} \in \mathbf{N} \setminus \{1\}, \quad n = 1, 2, 3, \dots \quad (2)$$

Initially, the potential is defined on the cube Q_0 . We want to extend it Q_0 -periodically to the entire lattice \mathbf{Z}^d , but we can only do that in steps: first extend it Q_0 -periodically to the cube Q_1 ; then extend the resulting potential Q_1 -periodically to the cube Q_2 ; etc. Assume, furthermore, that the process of periodic extension of the potential from the cube Q_{n-1} to the cube Q_n is imperfect: the values of the resulting Q_{n-1} -periodic potential on Q_n are then changed arbitrarily at m_n points chosen at random so that all such choices are equiprobable.

The potential v is, therefore, “grown” by the alternating process of periodic extension and introduction of “impurities”. The Borel-Cantelli lemma guarantees that $v(\cdot)$ is almost surely well-defined if

$$\sum_{n=1}^{\infty} \frac{m_n}{T_n^d} < \infty. \quad (3)$$

In general, the potential v is unbounded.

The goal of this work is to show that if the sequence T_n grows fast enough while the sequence m_n grows not too fast, then almost surely the operator h_v does not have eigenvalues – in other words, there are no localized states.¹ Before formulating the result precisely, we will describe our model in an equivalent but more convenient way.

We consider a Schrödinger operator in $l^2(\mathbf{Z}^d)$

$$h_v = \Delta + v, \quad (4)$$

where Δ is the lattice Laplacian:

$$(\Delta y)(x) = \sum_{\|z\|_1=1} y(x+z), \quad x \in \mathbf{Z}^d,$$

and v is the operator of multiplication by the function

$$v(x) = \sum_{n=0}^{\infty} u_n(x). \quad (5)$$

Here $u_n(\cdot)$ is a (T_n, \dots, T_n) -periodic function on \mathbf{Z}^d , where the numbers $T_n \in \mathbf{N}$ ($n = 0, 1, 2, \dots$) satisfy the condition (2). We assume that the function $u_n(\cdot)$ ($n \geq 1$) on its elementary cell Q_n defined by (1) vanishes everywhere except for a subset X_n of Q_n which consists of m_n points selected at random so that all such subsets of Q_n are equiprobable. The values of $u_n(\cdot)$ at those points are chosen arbitrarily. Under the

¹Note that if $m_n = 0$ for $n \geq n_0$, then the potential v is periodic; the spectrum of h_v in this case is well known to be not only continuous (no eigenvalues) but absolutely continuous.

condition (3), the potential v is almost surely well-defined, and the operator h_v is an (in general, unbounded) operator in $l^2(\mathbf{Z}^d)$ with the domain

$$D(h_v) = \{y \in l^2(\mathbf{Z}^d): \sum_{x \in \mathbf{Z}^d} |v(x)y(x)|^2 < \infty\}.$$

In our model, the functions u_n (and hence their sum v) are real-valued so the operator h_v is self-adjoint. However, in what follows the functions u_n may be complex-valued.

Let

$$W_n = T_n^{2d-1} \ln^2 T_n$$

($\ln^2 x$ means $(\ln x)^2$).

Theorem 1 *If*

$$\sum_{n=0}^{\infty} \frac{m_{n+1} W_n}{T_{n+1}} < \infty, \tag{6}$$

then with probability 1 the operator (4), (5) has no eigenvalues.

Here is a more detailed formulation of Theorem 1. By (Ω, \mathcal{F}, P) we denote the probability space of the model. Let the numbers $T_n \in \mathbf{N}$ ($n \geq 0$) satisfy the condition (2), and $X_n(\omega)$, for each $n \geq 1$, be a uniformly distributed random subset of cardinality m_n of the cube (1). Under condition (6), there is a set $\Omega_0 \in \mathcal{F}$ with $P(\Omega_0) = 1$ and the following property: suppose $\omega \in \Omega_0$ and $u_0(\cdot), u_1(\cdot), u_2(\cdot), \dots$ are arbitrary complex-valued functions such that each $u_n(\cdot)$ is (T_n, \dots, T_n) -periodic and, if $n \geq 1$, vanishes on the set $Q_n \setminus X_n(\omega)$; then the operator (4), (5) does not have eigenvalues.

The proof of Theorem 1 is based on a theorem regarding solutions of periodic lattice equations. Before formulating it, we introduce some notation. For a function $y: \mathbf{Z}^d \rightarrow \mathbf{C}$ let

$$\sigma(k) = \sum_{x \in \mathbf{Z}^d: x_d = k} |y(x)|^2, \quad k \in \mathbf{Z} \quad (0 \leq \sigma(k) \leq \infty).$$

Furthermore, we will represent any $T = \langle T_1, \dots, T_d \rangle \in \mathbf{N}^d$ in the form $T = \langle J, K \rangle$, where $J = \langle J_1, \dots, J_{d-1} \rangle \in \mathbf{N}^{d-1}$ and $K \in \mathbf{N}$ (so that $J_i = T_i$, $i = 1, \dots, d-1$, and $K = T_d$). The product $J_1 \dots J_{d-1}$ will be denoted by $[J]$.

Theorem 2 *For any $c > 0$, there is $\gamma_c > 0$ such that if $y: \mathbf{Z}^d \rightarrow \mathbf{C}$ is a solution of $(\Delta + v)y = \lambda y$, where $v(\cdot)$ is a (T_1, \dots, T_d) -periodic complex-valued function on \mathbf{Z}^d , and*

$$\sigma(k) < \infty \quad \text{for all } k \in \mathbf{Z}, \tag{7}$$

then

$$\sum_{1 \leq |r| \leq \rho_c(J)} \sigma(rK) \geq \gamma_c \sigma(0), \quad (8)$$

where

$$\rho_c(J) = (c[J]^2 \ln^2[J]) \vee (2[J]). \quad (9)$$

Remark. The inequalities (7) hold if and only if $\sigma(k_0) + \sigma(k_0 + 1) < \infty$ for some $k_0 \in \mathbf{Z}$. (This follows from the inequalities (34) and (35) below.)

Corollary 1 *For any $c > 0$, there is $\gamma_c > 0$ such that if $v(\cdot)$ is a (T_1, \dots, T_d) -periodic function on \mathbf{Z}^d and $y: \mathbf{Z}^d \rightarrow \mathbf{C}$ is a solution of $(\Delta + v)y = \lambda y$, then*

$$\sum_{K \leq |k| \leq \rho_c(J)K} \sigma(k) \geq \gamma_c \sigma(0), \quad (10)$$

where $\rho_c(J)$ is defined by (9).

The proof of Theorem 2 is based on a lower bound for iterates of a finite-dimensional linear operator.

Theorem 3 *Let A be an invertible linear operator in a finite-dimensional complex vector space X endowed with a seminorm $\|\cdot\|$. Then for any $g > 0$ and any $\xi \in X$*

$$\sum_{1 \leq |l| \leq (gn^2 \ln^2 n) \vee n} \|A^l \xi\|^2 \geq \mu_g \|\xi\|^2,$$

where $n = \dim X$ and $\mu_g > 0$ is a constant depending only on g .

2 Proofs

2.1 Proof of Theorem 3

Our starting point is the following result that has been useful in the theory of one-dimensional Schrödinger operators.

Lemma 1 ² *Let A be an invertible linear operator in an n -dimensional complex vector space X , $\|\cdot\|$ any seminorm on X , and ξ any vector in X . Then*

$$\max_{j=\pm 1, \dots, \pm n} \|A^j \xi\| \geq \frac{1}{n} \|\xi\|. \quad (11)$$

²This fact was first noticed, proved and used in [2] in the case where A is a unimodular operator in a two-dimensional real vector space. That was popularized by B.Simon as ‘‘Gordon’s lemma’’. However, the proof that appeared in B.Simon’s textbooks, as well as the general case of the lemma, actually belongs to E.A.Gorin (unpublished) and was communicated to B.Simon by S.Molchanov during B.Simon’s visit to Moscow in 1981 (without mentioning Gorin’s name).

Proof (E.A.Gorin). Let $p(x) = \sum_{k=0}^n b_k x^k$ be the characteristic polynomial of the operator A , and k_0 ($0 \leq k_0 \leq n$) be such that $|b_{k_0}| = \max_{0 \leq k \leq n} |b_k|$. By the Cayley-Hamilton theorem, $p(A)\xi = \sum_{k=0}^n b_k A^k \xi = 0$, so that

$$\xi = -b_{k_0}^{-1} A^{-k_0} \sum_{0 \leq k \leq n; k \neq k_0} b_k A^k \xi = - \sum_{0 \leq k \leq n; k \neq k_0} (b_k / b_{k_0}) A^{k-k_0} \xi,$$

and

$$\|\xi\| \leq \sum_{-k_0 \leq l \leq n-k_0; l \neq 0} \|A^l \xi\|, \quad (12)$$

which implies (11). ■

Now we will prove Theorem 3. In view of (12), we have:

$$\sum_{1 \leq |l| \leq n} \|A^l \xi\| \geq \|\xi\|. \quad (13)$$

Denote $\|A^m \xi\|$ by a_m for all $m \in \mathbf{Z}$. We may assume that $a_0 = 1$. Pick an arbitrary integer $k \geq 1$ (its actual value will be selected later). It follows from the inequality (13) applied to A^j instead of A that

$$\sum'_{l=-n}^n a_{l \cdot j} \geq 1, \quad j = 1, 2, \dots, k$$

(a prime following the summation symbol denotes the omission of the zero value of the summation index). For any $s \in \mathbf{Z}$ ($s \neq 0$), let d_s be the number of distinct positive divisors of s (including 1 and $|s|$). The previous inequality and the Cauchy – Schwartz inequality imply:

$$\sum'_{l=-n}^n \frac{a_{l \cdot j}^2}{d_{l \cdot j}} \cdot \sum'_{l=-n}^n d_{l \cdot j} \geq 1,$$

or

$$\sum'_{l=-n}^n \frac{a_{l \cdot j}^2}{d_{l \cdot j}} \geq \left(\sum'_{l=-n}^n d_{l \cdot j} \right)^{-1}, \quad j = 1, 2, \dots, k.$$

Summation on j , $1 \leq j \leq k$, gives

$$\sum'_{s=-kn}^{kn} \frac{a_s^2}{d_s} \cdot \Phi(s) \geq \sum_{j=1}^k \left(\sum'_{l=-n}^n d_{l \cdot j} \right)^{-1}, \quad (14)$$

where $\Phi(s)$ is the number of representations of s in the form $s = l \cdot j$ with $-n \leq l \leq n$, $1 \leq j \leq k$. Obviously, $\Phi(s) \leq d_s$, so that (14) implies

$$\sum'_{s=-kn}^{kn} a_s^2 \geq \sum_{j=1}^k \left(\sum'_{l=-n}^n d_{l \cdot j} \right)^{-1}. \quad (15)$$

On the other hand, applying the inequality between the arithmetic mean and the harmonic mean

$$\frac{1}{k} \sum_{j=1}^k x_j \geq \left(\frac{1}{k} \sum_{j=1}^k x_j^{-1} \right)^{-1} \quad (\text{all } x_j > 0),$$

using then the obvious relation $d_{l,j} \leq d_l \cdot d_j$ and putting, for any positive integer p ,

$$D_p = \sum_{s=1}^p d_s, \quad (16)$$

we get

$$\sum_{j=1}^k \left(\sum_{l=-n}^n d_{l,j} \right)^{-1} \geq k^2 \left(\sum_{j=1}^k \sum_{l=-n}^n d_{l,j} \right)^{-1} \geq k^2 \left(\sum_{j=1}^k \sum_{l=-n}^n d_l \cdot d_j \right)^{-1} = k^2 (2D_k D_n)^{-1}. \quad (17)$$

Note that

$$D_p = \sum_{i=1}^p \left\lfloor \frac{p}{i} \right\rfloor,$$

where $\lfloor x \rfloor$ is the integer part of x . (Proof: rewrite (16) in the form $D_p = \sum_{s=1}^p \sum_{i=1}^p y_{is}$, where y_{is} equals 1 if i is a divisor of s and equals 0 otherwise; then reverse the order of summation. For a different proof of this simple fact, see [5, Problem VIII 79].) It follows that

$$D_p \leq p \sum_{i=1}^p \frac{1}{i} \leq p(\ln p + 1),$$

so that (15) and (17) imply:

$$\sum_{s=-kn}^{kn} a_s^2 \geq \frac{k^2}{2kn(\ln k + 1)(\ln n + 1)} = \frac{k}{2n(\ln k + 1)(\ln n + 1)}. \quad (18)$$

This is true for all integers $k \geq 1$. Now we put $k = \lfloor gn \ln^2 n \vee 1 \rfloor$, where $g > 0$. The right hand side of (18) (now depending only on n) is positive for all $n \geq 1$; as $n \rightarrow \infty$, it approaches $g/2 > 0$. Therefore, it exceeds some number $\mu_g > 0$ for all $n \geq 1$. This proves Theorem 3. ■

2.2 Proof of Theorem 2

(a) Let $y(\cdot)$ be a solution of the equation $(\Delta + v)y = \lambda y$ with a (J, K) -periodic $v(\cdot)$ such that $\sigma(k) < \infty$ for all $k \in \mathbf{Z}$. By including λ in v we may assume that $\lambda = 0$ so that

$$(\Delta + v)y = 0. \quad (19)$$

We will use the representation of any $x \in \mathbf{Z}^d$ in the form $x = \langle j, k \rangle$, where $j = \langle j_1, \dots, j_{d-1} \rangle \in \mathbf{Z}^{d-1}$ and $k \in \mathbf{Z}$; then (19) takes the form

$$y(j, k+1) + y(j, k-1) + \sum_{\substack{1 \leq s \leq d-1 \\ \varepsilon = \pm 1}} y(j + \varepsilon e_s, k) + v(j, k)y(j, k) = 0, \quad j \in \mathbf{Z}^{d-1}, \quad k \in \mathbf{Z}, \quad (20)$$

where $e_s = \langle 0, \dots, 0, 1, 0, \dots, 0 \rangle \in \mathbf{Z}^{d-1}$ (1 is in the s th place).

(b) For $\theta = \langle \theta_1, \dots, \theta_{d-1} \rangle \in \mathbf{T}^{d-1}$ let

$$z(j, k; \theta) = y(j, k)e^{2\pi i(j \cdot \theta)}, \quad j \in \mathbf{Z}^{d-1}, \quad k \in \mathbf{Z},$$

where $i = \sqrt{-1}$ and

$$j \cdot \theta = j_1 \theta_1 + \dots + j_{d-1} \theta_{d-1}.$$

The equation (20) implies that

$$z(j, k+1) + z(j, k-1) + \sum_{\substack{1 \leq s \leq d-1 \\ \varepsilon = \pm 1}} e^{-2\pi i \varepsilon \theta_s} z(j + \varepsilon e_s, k) + v(j, k)z(j, k) = 0, \quad j \in \mathbf{Z}^{d-1}, \quad k \in \mathbf{Z}. \quad (21)$$

(c) Consider the “discrete cylinder” C_J :

$$C_J = \mathbf{Z}_J^{d-1} \times \mathbf{Z},$$

where

$$\mathbf{Z}_J^{d-1} = (\mathbf{Z}/J_1\mathbf{Z}) \times \dots \times (\mathbf{Z}/J_{d-1}\mathbf{Z}) \simeq \mathbf{Z}^{d-1}/G_J.$$

Here G_J is the following subgroup of \mathbf{Z}^{d-1} :

$$G_J = \{x = \langle x_1, \dots, x_{d-1} \rangle \in \mathbf{Z}^{d-1}: x_1 \in J_1\mathbf{Z}, \dots, x_{d-1} \in J_{d-1}\mathbf{Z}\}.$$

The function $v(j, k)$, being (J, K) -periodic, carries over in a natural manner to C_J , giving rise to a function $V(\nu, k)$ ($\nu \in \mathbf{Z}_J^{d-1}$, $k \in \mathbf{Z}$), which is K -periodic in k :

$$V(\nu, k+K) = V(\nu, k), \quad \nu \in \mathbf{Z}_J^{d-1}, \quad k \in \mathbf{Z}. \quad (22)$$

(d) For each pair $\langle \nu, k \rangle \in C_J$, define a function $Z(\nu, k; \cdot) \in L^2(\mathbf{T}^{d-1})$ as the sum of the series

$$Z(\nu, k; \theta) = \sum_{j \in \nu} z(j, k; \theta) \equiv \sum_{j \in \nu} y(j, k)e^{2\pi i(j \cdot \theta)}, \quad (23)$$

which converges in $L^2(\mathbf{T}^{d-1})$ due to (7).

Fix $\langle \nu, k \rangle \in C_J$. Summation on $j \in \nu$ in (21) gives the following equation in $L^2(\mathbf{T}^{d-1})$:

$$Z(\nu, k+1; \theta) + Z(\nu, k-1; \theta) + \sum_{\substack{1 \leq s \leq d-1 \\ \varepsilon = \pm 1}} e^{-2\pi i \varepsilon \theta_s} Z(\nu + \varepsilon e_s, k; \theta) + V(\nu, k) Z(\nu, k; \theta) = 0. \quad (24)$$

Each $Z(\nu, k; \cdot)$, as an element of $L^2(\mathbf{T}^{d-1})$, is a class of equivalent square integrable Lebesgue measurable functions on \mathbf{T}^{d-1} ; select a Borel measurable function in this class. From now on $Z(\nu, k; \cdot)$ will denote that function on \mathbf{T}^{d-1} , and each equation (24) will be considered as a pointwise equation that holds for Lebesgue almost every $\theta \in \mathbf{T}^{d-1}$. Let B be a full measure Borel subset of \mathbf{T}^{d-1} such that for any $\theta \in B$ all the equations (24) hold.

(e) Pick an arbitrary $\theta \in B$. For any $k \in \mathbf{Z}$, consider a vector

$$\zeta_k(\theta) = \langle \mathbf{Z}(\nu, k; \theta) \rangle_{\nu \in \mathbf{Z}_J^{d-1}} \in \mathcal{M},$$

where \mathcal{M} is the $[J]$ -dimensional coordinate vector space over \mathbf{C} with coordinates labeled by elements of \mathbf{Z}_J^{d-1} . The equation (24) can be rewritten in the form

$$\zeta_{k+1}(\theta) = -\zeta_{k-1}(\theta) + L_k(\theta)\zeta_k(\theta), \quad (25)$$

where $L_k(\theta)$ is the linear operator in \mathcal{M} acting on vectors $\langle \zeta(\nu) \rangle_{\nu \in \mathbf{Z}_J^{d-1}} \in \mathcal{M}$ as follows:

$$(L_k(\theta)\zeta)(\nu) = - \sum_{\substack{1 \leq s \leq d-1 \\ \varepsilon = \pm 1}} e^{-2\pi i \varepsilon \theta_s} \zeta(\nu + \varepsilon e_s) - V(\nu, k)\zeta(\nu).$$

In view of (22), for any $k \in K$

$$L_{k+K}(\theta) = L_k(\theta).$$

(f) Let $\mathcal{L} = \mathcal{M} \oplus \mathcal{M}$ and

$$\xi_k(\theta) = \langle \zeta_k(\theta), \zeta_{k-1}(\theta) \rangle \in \mathcal{L}.$$

The equation (25) can be further transformed into

$$\xi_{k+1}(\theta) = T_k(\theta)\xi_k(\theta),$$

where $T_k(\theta)$ is the following linear operator in \mathcal{L} :

$$T_k(\theta)(\langle \zeta, \eta \rangle) = \langle -\eta + L_k(\theta)\zeta, \zeta \rangle.$$

Again, $T_k(\theta)$ is K -periodic in k :

$$T_{k+K}(\theta) = T_k(\theta). \quad (26)$$

Let, furthermore,

$$S(\theta) = T_{K-1}(\theta)T_{K-2}(\theta) \dots T_1(\theta)T_0(\theta).$$

By (26), we have

$$\xi_{rK}(\theta) = (S(\theta))^r \xi_0(\theta) \quad (27)$$

for all $r \in \mathbf{Z}$.

(g) Let $\|\cdot\|$ be an arbitrary seminorm in \mathcal{L} . Put $g = c/16$, where c is the positive parameter in Theorem 2. By (27) and Theorem 3,

$$\sum_{1 \leq |r| \leq \beta(J)} \|\xi_{rK}(\theta)\|^2 \geq \mu_g \|\xi_0(\theta)\|^2, \quad (28)$$

where

$$\beta(J) = g(2[J] \ln(2[J]))^2 \vee (2[J]).$$

Note that if $[J] > 1$ then $\beta(J) \leq c([J] \ln[J])^2 \vee (2[J]) = \rho_c(J)$, so by (28) we have:

$$\sum_{1 \leq |r| \leq \rho_c(J)} \|\xi_{rK}(\theta)\|^2 \geq \mu_g \|\xi_0(\theta)\|^2.$$

If $[J] = 1$, the inequality $\rho_c(J) \geq 2$ implies that

$$\sum_{1 \leq |r| \leq \rho_c(J)} \|\xi_{rK}(\theta)\|^2 \geq \sum_{1 \leq |r| \leq 2} \|\xi_{rK}(\theta)\|^2 \geq \frac{1}{4} \|\xi_0(\theta)\|^2$$

(the second inequality follows from Lemma 1 with $n = 2$). Therefore, in all cases

$$\sum_{1 \leq |r| \leq \rho_c(J)} \|\xi_{rK}(\theta)\|^2 \geq \gamma_c \|\xi_0(\theta)\|^2, \quad (29)$$

where $\gamma_c = (\mu_{c/16}) \wedge (1/4)$.

(h) Choose the following seminorm in \mathcal{L} :

$$\|\langle \zeta, \eta \rangle\| = \left(\sum_{\nu \in \mathbf{Z}_J^{d-1}} |\zeta_\nu|^2 \right)^{1/2}. \quad (30)$$

Then (29) becomes

$$\sum_{1 \leq |r| \leq \rho_c(J)} \left(\sum_{\nu \in \mathbf{Z}_J^{d-1}} |Z(\nu, rK; \theta)|^2 \right) \geq \gamma_c \sum_{\nu \in \mathbf{Z}_J^{d-1}} |Z(\nu, 0; \theta)|^2, \quad (31)$$

which is true for all $\theta \in B$.

Now integrate both sides of this inequality over B or, equivalently, over \mathbf{T}^{d-1} . By Parseval's identity, (23) implies that

$$\int_{\mathbf{T}^{d-1}} |Z(\nu, k, \theta)|^2 d\theta = \sum_{j \in \nu} |y(j, k)|^2,$$

so that (31) upon integration gives

$$\sum_{1 \leq |r| \leq \rho_c(J)} \left(\sum_{j \in \mathbf{Z}^{d-1}} |y(j, rK)|^2 \right) \geq \gamma_c \sum_{j \in \mathbf{Z}^{d-1}} |y(j, 0)|^2,$$

or

$$\sum_{1 \leq |r| \leq \rho_c(J)} \sigma(rK) \geq \gamma_c \sigma(0).$$

This completes the proof of Theorem 2. ■

Remark. The following version of Theorem 2 is true without the assumption (7).

Theorem 4 *If $y: \mathbf{Z}^d \rightarrow \mathbf{C}$ is a solution of $(\Delta + v)y = \lambda y$, where $v(\cdot)$ is (T_1, \dots, T_d) -periodic on \mathbf{Z}^d , then for any $c > 0$*

$$\sum_{1 \leq |r| \leq \rho_c(J)} (\sigma(rK) + \sigma(rK - 1)) \geq \gamma_c (\sigma(0) + \sigma(-1)), \quad (32)$$

where $\rho_c(J)$ and γ_c are the same as in Theorem 2.

Proof. We may assume that the left-hand side of (32) is finite. Then

$$\sigma(K) + \sigma(K - 1) < \infty. \quad (33)$$

Let $y_k = \langle y(j, k) \rangle_{j \in \mathbf{Z}^{d-1}}$ so that $\|y_k\|_2 = (\sigma(k))^{1/2}$. In view of (20) and Minkowski's inequality,

$$(\sigma(k + 1))^{1/2} \leq (\sigma(k - 1))^{1/2} + (2(d - 1) + \|v\|_\infty)(\sigma(k))^{1/2} \quad (34)$$

and similarly

$$(\sigma(k-1))^{1/2} \leq (\sigma(k+1))^{1/2} + (2(d-1) + \|v\|_\infty)(\sigma(k))^{1/2}, \quad (35)$$

so (33) implies that $\sigma(k) < \infty$ for all $k \in \mathbf{Z}$. The remaining part of the proof of Theorem 4 is identical to the proof of Theorem 2 except for the last step when we choose, instead of (30), the standard Euclidean norm in \mathcal{L} :

$$\|\langle \zeta, \eta \rangle\| = \left(\sum_{\nu \in \mathbf{Z}_J^{d-1}} |\zeta_\nu|^2 + \sum_{\nu \in \mathbf{Z}_J^{d-1}} |\eta_\nu|^2 \right)^{1/2}$$

and, instead of (8), obtain (32). ■

Proof of Corollary 1

We may assume that the left-hand side of (10) is finite. Since $\rho_c(J) \geq 2$, we have $\sigma(K) + \sigma(K+1) < \infty$ and hence, in view of the inequalities (34) and (35), $\sigma(k) < \infty$ for all $k \in \mathbf{Z}$. Now (10) follows from Theorem 2. ■

2.3 Proof of Theorem 1

For all $n \geq 0$ and $x \in \mathbf{Z}^d$, let

$$v_n(x) = u_0(x) + \dots + u_n(x).$$

Consider the elementary cell of $v_{n+1}(\cdot)$:

$$Q_{n+1} = \{x \in \mathbf{Z}^d: -T_{n+1}/2 < x_j \leq T_{n+1}/2, \quad j = 1, \dots, d\}$$

and the set

$$L_n = \{x \in \mathbf{Z}^d: -W_n \leq x_d \leq W_n\}.$$

Let

$$w_{n+1} = \Pr\{X_{n+1} \cap L_n \neq \emptyset\} \equiv \Pr\{X_{n+1} \cap R_{n+1} \neq \emptyset\},$$

where $R_{n+1} = Q_{n+1} \cap L_n$. We have

$$w_{n+1} \leq m_{n+1} \frac{2W_n + 1}{T_{n+1}}. \quad (36)$$

Indeed, for any point $x \in Q_{n+1}$ the probability h_{n+1} of the event $x \in X_{n+1}$ is m_{n+1}/T_{n+1}^d ; therefore, (36) is a consequence of the following inequalities (in which \mathbf{E} denotes the expectation):

$$w_{n+1} \leq \mathbf{E}[\text{Card}(X_{n+1} \cap R_{n+1})] = \text{Card}(R_{n+1})h_{n+1} \leq (2W_n + 1)T_{n+1}^{d-1} h_{n+1}.$$

By the Borel-Cantelli Lemma, (6) implies that with probability 1 there is a random number $N(\omega)$ such that $X_{n+1} \cap L_n = \emptyset$ for all $n \geq N(\omega)$. Consequently, for all $n \geq N(\omega)$ we have $(X_{n+1} \cup X_{n+2} \cup \dots) \cap L_n = \emptyset$ so that on the set L_n $u_{n+1}(x) = u_{n+2}(x) = \dots = 0$ and, therefore,

$$v(x) = v_n(x) \quad \text{on } L_n \text{ for } n \geq N(\omega). \quad (37)$$

Suppose $y(\cdot)$ is a solution of the equation $(\Delta + v)y = \lambda y$ that belongs to $l^2(\mathbf{Z}^d)$ so that

$$\sum_{k \in \mathbf{Z}} \sigma(k) < \infty.$$

Fix temporarily some $n \geq 0$ and denote by $y_n(\cdot)$ the function on \mathbf{Z}^d satisfying the equation

$$(\Delta + v_n)y_n = \lambda y_n$$

and such that $y_n(j, k) = y(j, k)$ for $k = -1, 0$ and all $j \in \mathbf{Z}^{d-1}$. Such function y_n exists uniquely; put

$$\sigma_n(k) = \sum_{j \in \mathbf{Z}^{d-1}} |y_n(j, k)|^2, \quad k \in \mathbf{Z}.$$

By Corollary 1 to Theorem 2 (with $J = \langle T_n, \dots, T_n \rangle \in \mathbf{N}^{d-1}$ and $K = T_n$), for any $c > 0$

$$\sum_{T_n \leq |k| \leq q_c(n)} \sigma_n(rT_n) \geq \gamma_c \sigma_n(0), \quad (38)$$

where

$$q_c(n) = \left(c(T_n^{2d-2} \ln^2 T_n^{d-1}) \vee (2T_n^{d-1}) \right) T_n$$

and $\gamma_c > 0$. Putting $c = (d-1)^{-2}$, for large enough n ($n \geq N_0$) we have $q_c(n) = T_n^{2d-1} \ln^2 T_n = W_n$ so that (38) becomes

$$\sum_{T_n \leq |k| \leq W_n} \sigma_n(k) \geq \gamma \sigma_n(0), \quad (39)$$

where $\gamma = \gamma_{(d-1)^{-2}}$.

Due to (37), we have $y_n(x) = y(x)$ on the set L_n if $n \geq N(\omega)$, and hence

$$\sum_{T_n \leq |k| \leq W_n} \sigma(k) \geq \gamma \sigma(0)$$

if $n \geq N(\omega) \vee N_0$. Therefore, almost surely (a.s.)³

$$\sigma(0) \leq \gamma^{-1} \sum_{|k| \geq T_n} \sigma(k) \quad (40)$$

for all large enough n ; since $\sum_{k \in \mathbf{Z}} \sigma(k) < \infty$ and $T_n \rightarrow \infty$ as $n \rightarrow \infty$, it follows from (40) that $\sigma(0) = 0$, that is, $y(j, 0) = 0$ for all $j \in \mathbf{Z}^{d-1}$.

Similarly, the function $y^*(j, k) = y(j, k + 1)$ satisfies the equation $(\Delta + v^*)y^* = \lambda y^*$, where $v^*(j, k) = v(j, k + 1)$ for all j, k . The potential v^* has the same structure as v , so the inclusion $y^* \in l^2(\mathbf{Z}^d)$ a.s. implies that $y^*(j, 0) = 0$ for all j or, equivalently, $y(j, 1) = 0$ for all j .

Therefore, if $(\Delta + v)y = \lambda y$ and $y \in l^2(\mathbf{Z}^d)$, then a.s. $y(j, 0) = y(j, 1) = 0$ for all $j \in \mathbf{Z}^{d-1}$ and hence $y \equiv 0$ on \mathbf{Z}^d . This completes the proof of Theorem 1. ■

3 Concluding Remarks

1. The transformation of multidimensional lattice equations close to periodic ones into quasi-1D equations, in combination with lower bounds for solutions of periodic quasi-1D equations, was first used in [3], where the absence of eigenvalues for a lattice operator with a fast uniform periodic approximation was proved. This idea was then used in [1] in a different situation, where the potential is periodic in $d - 1$ coordinate directions and admits a fast local periodic approximation in the remaining direction.

2. Since in the model considered above the nonzero values of $u_n(\cdot)$ are arbitrary, the potential $v(x)$ in Theorem 1 can grow (with respect to the distance of x from the origin) arbitrarily fast on a suitable random subset of the lattice; moreover, by choosing appropriate (say, nonnegative) values of the summands $u_n(\cdot)$ we can make the average

$$A(r) = \frac{\sum_{x \in B_r} v(x)}{|B_r|}$$

of v over the cube $B_r = \{x \in \mathbf{Z}^d \mid \|x\|_\infty \leq r\}$ grow arbitrarily fast as $r \in \mathbf{N}$ goes to ∞ ; the operator (4) will still have no eigenvalues a.s.

³i.e., for P -a.e. ω and any $v = u_0 + u_1 + u_2 + \dots$, where each u_n ($n \geq 1$) vanishes on $Q_n \setminus X_n(\omega)$.

3. In the model above, the intersection $X_n(\omega) \cap Q_{n-1}$ can be nonempty, i.e., the impurities introduced during the recursive process of the “growth” of the potential can affect its portions constructed earlier. This will not occur if we modify the definition of the random set $X_n(\omega)$ making it a uniformly distributed subset of cardinality m_n of the set $Q_n \setminus Q_{n-1}$. This change in the model does not affect the validity of Theorem 1 and requires almost no change in its proof.

4. The model can also be modified by assuming that the function $u_n(\cdot)$ representing the impurities introduced at the n th step ($n \geq 1$) is defined as follows: for any point x of the cube Q_n , $u_n(x)$ is a random variable such that $\Pr\{u_n(x) \neq 0\} \leq p_n$. The corresponding version of Theorem 1 states: *If $\sum_{n \geq 0} T_{n+1}^{d-1} W_n p_{n+1} < \infty$, then with probability 1 the operator (4), (5) has no eigenvalues.*

5. The elementary cell of the function $u_n(\cdot)$ does not have to be a cube – it can be a “rectangular box”

$$Q_n = \{x \in \mathbf{Z}^d: -\frac{T_{n,i}}{2} < x_i \leq \frac{T_{n,i}}{2}, \quad i = 1, \dots, d\},$$

where $T_{n,i} \in \mathbf{N}$ and $T_{n,i}/T_{n-1,i} \in \mathbf{N} \setminus \{1\}$ ($i = 1, \dots, d$; $n \geq 1$). Let

$$T_{n,\hat{i}} = T_{n,1} \dots T_{n,i-1} T_{n,i+1} \dots T_{n,d}$$

and $W_{n,i} = T_{n,i}(T_{n,\hat{i}})^2 \ln^2 T_{n,\hat{i}}$. *If for some i ($1 \leq i \leq d$)*

$$\sum_{n=0}^{\infty} \frac{m_{n+1} W_{n,i}}{T_{n+1,i}} < \infty,$$

then with probability 1 the operator (4), (5) does not have eigenvalues.

6. It may seem that the absence of localized states in the model is caused by the fast decay, as n grows, of the “density” of the impurities introduced at the n th step. However, this is not the case. To see why, consider a deterministic version of the model, in which the (T_n, \dots, T_n) -periodic function $u_n(\cdot)$ vanishes on the set $L_{n-1} = \{x \in \mathbf{Z}^d: -W_{n-1} \leq x_d \leq W_{n-1}\}$ while its values on the set $Q_n \setminus L_{n-1}$ are arbitrary. The operator (4), (5) still has no eigenvalues.

7. The main results of the paper – Theorems 2 and 1 – pertain to the Schrödinger operator (4), but both theorems (and their proofs) can be adapted to other finite-range lattice operators, such as $\tilde{\Delta} + v$, where $\tilde{\Delta}$ is the Molchanov-Vainberg Laplacian (diagonal Laplacian) [4]:

$$(\tilde{\Delta}y)(x) = \sum_{z \in \{-1, +1\}^d} y(x+z), \quad x \in \mathbf{Z}^d.$$

8. Although the model considered in the paper is multi-dimensional, the “imperfectly grown” potentials can also be considered in dimension one. The absence of eigenvalues in this case is guaranteed by conditions (2) and (6) with $W_n = 2T_n$, and follows from Lemma 1. Combining this with the Simon-Spencer theorem [6] (in dimension one, potentials that are unbounded on both \mathbf{Z}_+ and \mathbf{Z}_- produce no absolutely continuous spectrum), we obtain yet another class of one-dimensional operators with purely singular continuous spectrum.

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References

- [1] D. Damanik, A version of Gordon’s theorem for multidimensional Schrödinger operators, *Trans. Amer. Math. Soc.* **356** (2004), 495–507 MR2022708.
- [2] A. Y. Gordon, On the point spectrum of one-dimensional Schrödinger operator, *Uspekhi Mat. Nauk (Russian Math. Surveys)* **31** (1976), 257–258 (in Russian) MR **56**:16450.
- [3] A. Y. Gordon, A sufficient condition for the continuity of the spectrum of a discrete Schrödinger operator, *Funct. Anal. Appl.* **20** (1986), 313–315 MR **88d**:39011.
- [4] S. Molchanov and B. Vainberg, Scattering on the system of the sparse bumps: multidimensional case, *Applicable Analysis* **71** (1999), 167–185 MR1690097 (2000a:47071).
- [5] G. Polya and G. Szegő, *Problems and theorems in analysis II*, Springer-Verlag, Berlin, 1976 MR 1492448.
- [6] B. Simon and T. Spencer, Trace class perturbations and the absence of absolutely continuous spectrum, *Communications in Mathematical Physics* **125** (1989), 113–125.