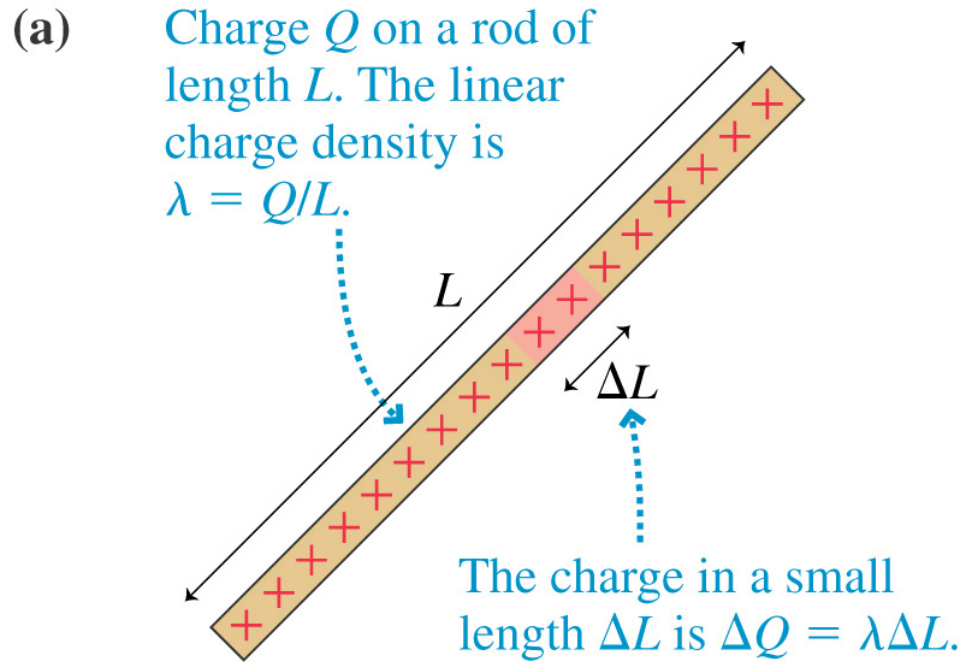


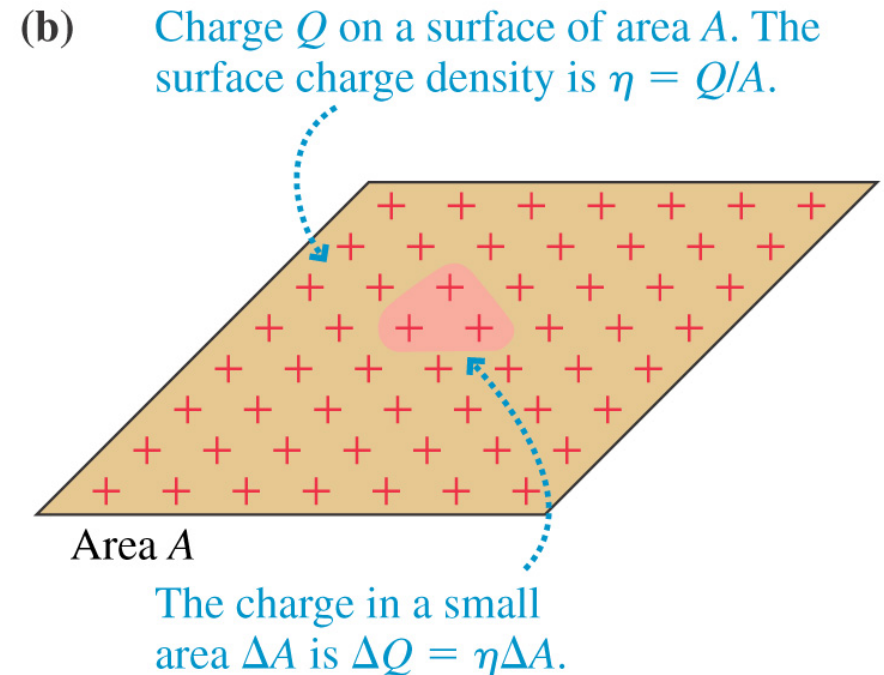
Lecture 4: Chapter 26, September 6 2005

The Electric Field of a Continuous Charge Distribution



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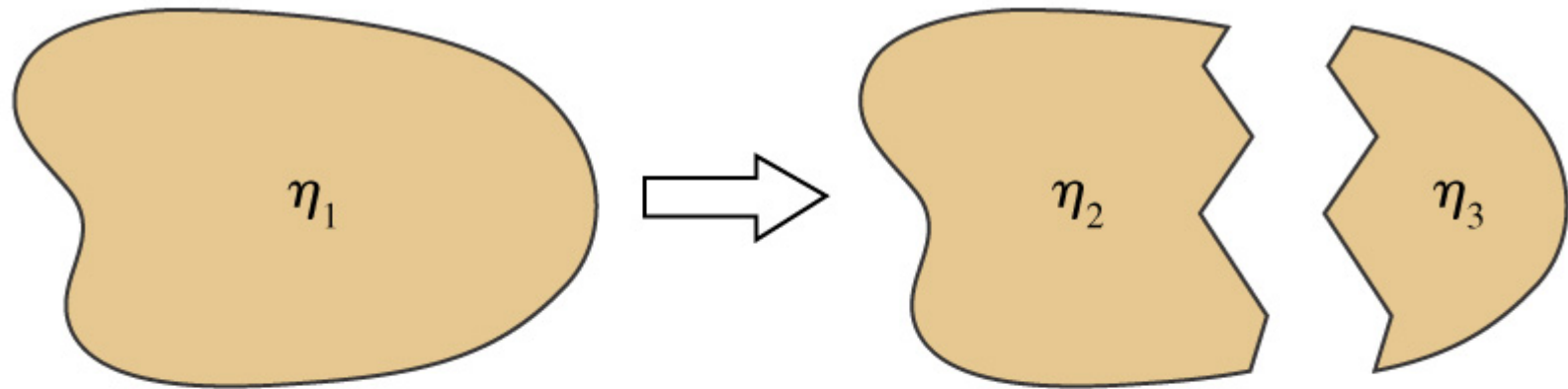
Linear charge density:
 $\lambda = Q/L$



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Surface charge density:
 $\eta = Q/A$ (In most of other texts it is σ)

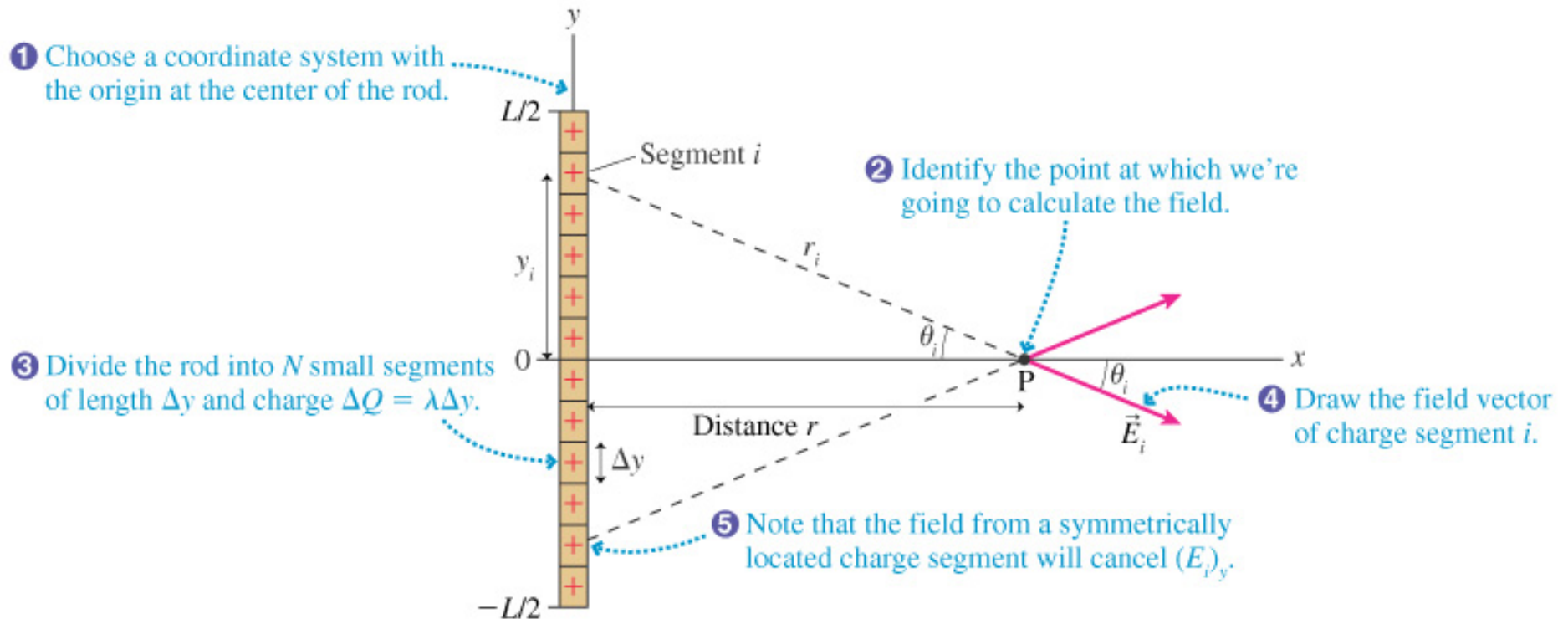
A piece of plastic is uniformly charged with surface density η_1 . After braking into pieces, rank in order, from largest to smallest, the surface charge densities η_1 to η_3 .



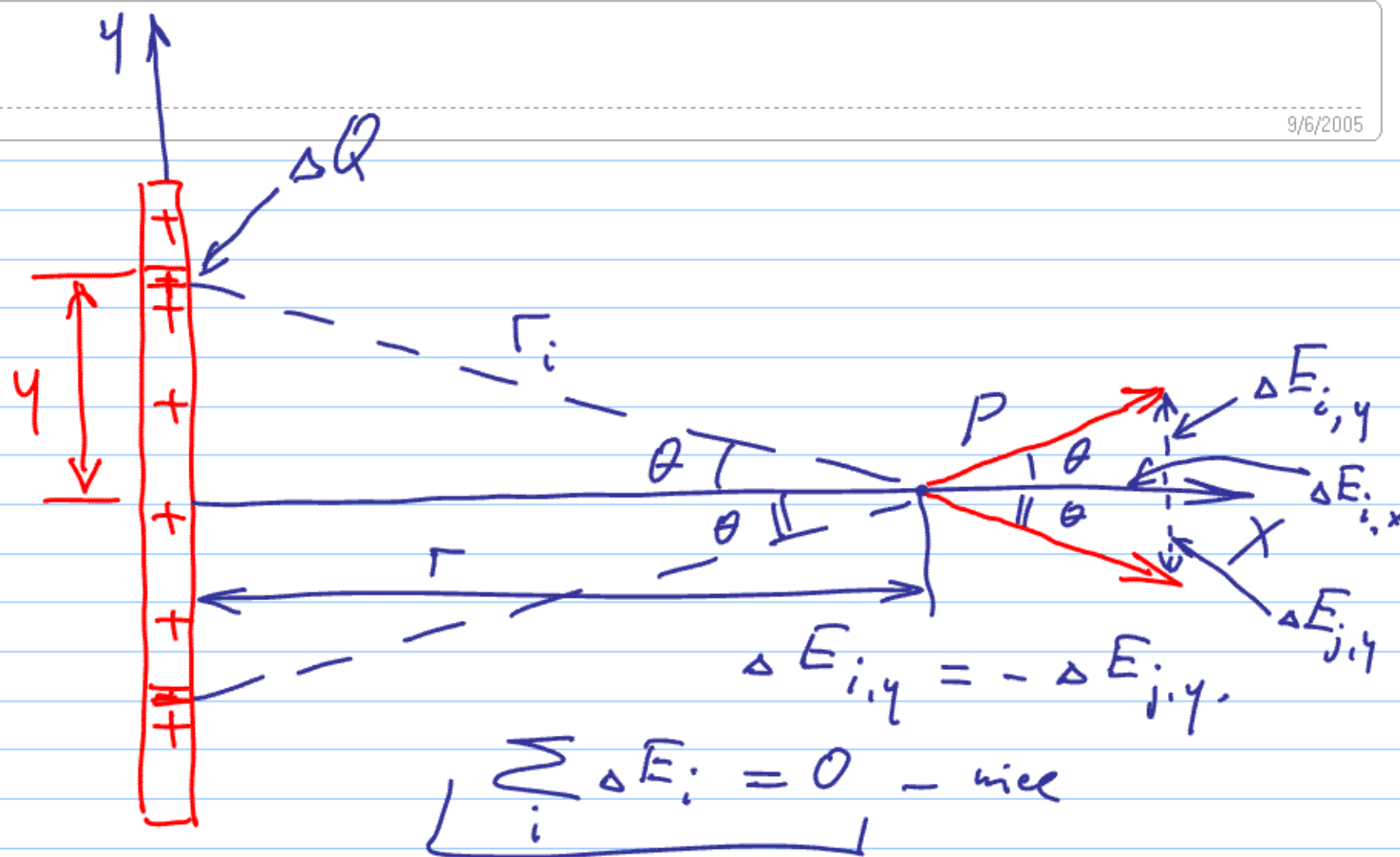
Is plastic dielectric or metal? Why is this important?

Problem Solving Strategy

Example: Electric field of a line of charge



- Divide the total charge Q into small pieces of charge ΔQ
- Use Coulomb law and draw the el. field ($\vec{E}_i$) for one or two small pieces of charge
- Look for symmetries of the charge distribution that simplify the fields
- Find $\vec{E}_{\text{net}} = \Sigma \vec{E}_i$, express an integral and integrate it component by component



$$\Delta E_{i,x} = |\Delta E_i| \cdot \cos \theta =$$

$$\begin{aligned}
 &= \frac{1}{4\pi\epsilon_0} \frac{\Delta Q_i \cdot \cos\theta}{r_i^2} = \frac{1}{4\pi\epsilon_0} \frac{\Delta Q_i}{r^2 + y_i^2} \left(\frac{r}{\sqrt{r^2 + y_i^2}} \right) \cos\theta \\
 &= \frac{1}{4\pi\epsilon_0} \frac{\Delta Q_i \cdot r}{(r^2 + y_i^2)^{3/2}}
 \end{aligned}$$

$$\begin{aligned}
 E_{\text{net}, x} &= \sum_i \Delta E_{i,x} = \frac{r}{4\pi\epsilon_0} \sum_i \frac{\Delta Q_i}{(r^2 + y_i^2)^{3/2}} \\
 \sum &\rightarrow \int \\
 &= \frac{r}{4\pi\epsilon_0} \int_{-\frac{L}{2}}^{\frac{L}{2}} \frac{\lambda \cdot dy}{(r^2 + y^2)^{3/2}} =
 \end{aligned}$$

$$\boxed{\int \frac{dy}{(r^2 + y^2)^{3/2}} = \frac{y}{r^2(r^2 + y^2)^{1/2}}}$$

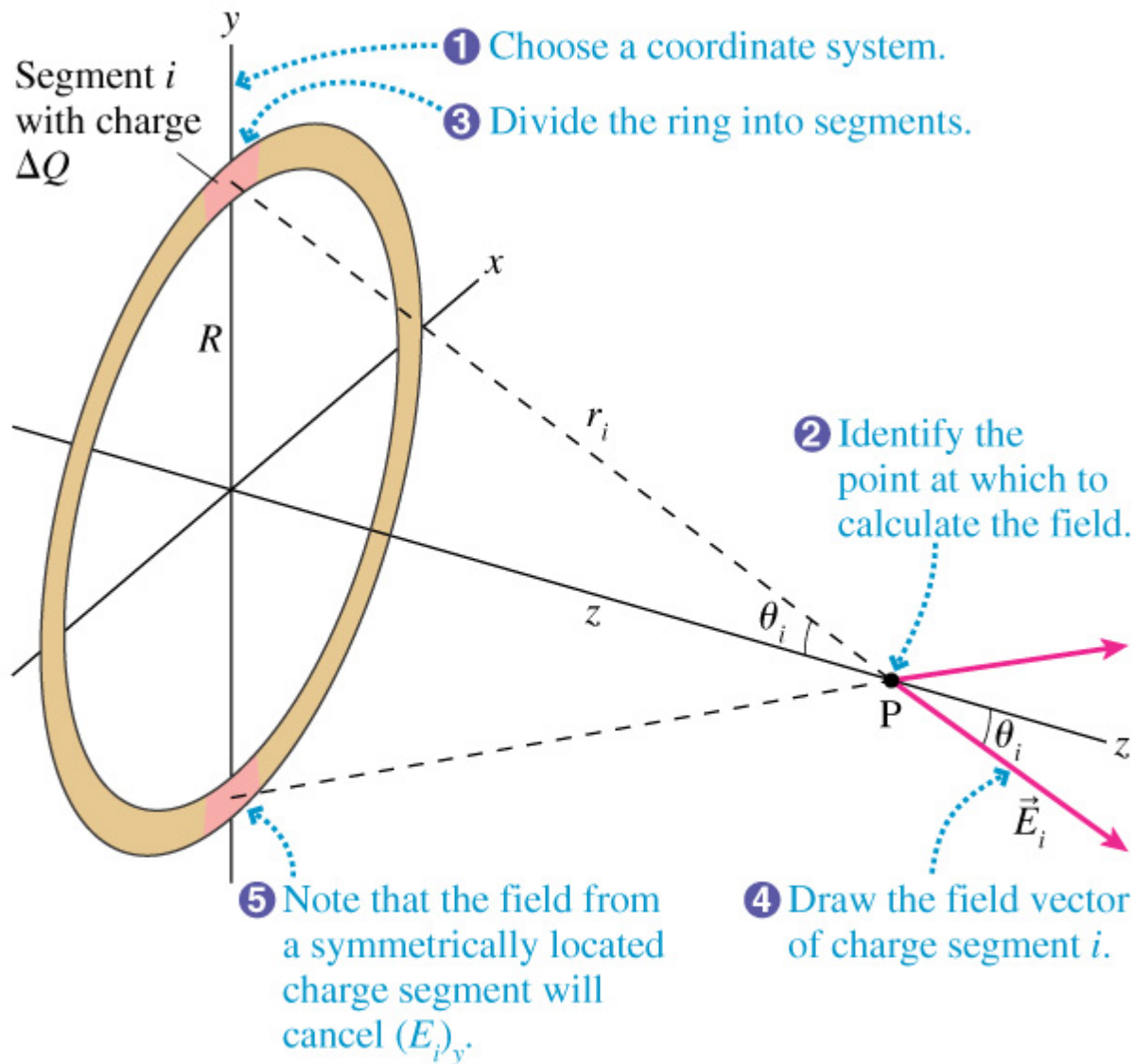
$$= \frac{\cancel{\lambda} \cdot \lambda}{4\pi\epsilon_0} \frac{y}{r^2 (r^2 + y^2)^{1/2}} \Big|_{-\frac{L}{2}}^{\frac{L}{2}} =$$

$$\frac{\lambda}{4\pi\epsilon_0 r} \left(\frac{L/2}{\sqrt{r^2 + (\frac{L}{2})^2}} - \frac{(-\frac{L}{2})}{\sqrt{r^2 + (\frac{L}{2})^2}} \right) =$$

$$= \frac{\lambda}{4\pi\epsilon_0 \cdot r} \frac{L}{\sqrt{r^2 + (\frac{L}{2})^2}} = \frac{\boxed{\lambda \cdot L = Q}}{4\pi\epsilon_0 \cdot r \cdot \sqrt{r^2 + (\frac{L}{2})^2}} \frac{Q}{1}$$

$$E_{\text{net}, x} (L \rightarrow \infty) = \frac{Q = \lambda \cdot L}{4\pi\epsilon_0} \frac{1}{r \frac{L}{2}} = \frac{2\lambda}{4\pi\epsilon_0 \cdot r}$$

The Electric Field of a Ring of Charge



See solution in Chapter 26, page 829

Distance dependencies of some distribution of charges

$$E_{\text{line}} = \lim_{L \rightarrow \infty} \frac{1}{4\pi\epsilon_0} \frac{|Q|}{r\sqrt{r^2 + (L/2)^2}} = \frac{1}{4\pi\epsilon_0} \frac{|Q|}{rL/2} = \frac{1}{4\pi\epsilon_0} \frac{2|\lambda|}{r}$$

An Infinite Line of Charge

- Line of charge is infinite in one dimension in 3-D space $\Rightarrow Q$ is infinite

$$E_{\text{point}} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2}$$

A point Charge

- *Point* charge is totally localized and Q (or q) is finite

$$\vec{E}_{\text{dipole}} \approx \frac{1}{4\pi\epsilon_0} \frac{2\vec{p}}{r^3} \quad (\text{on the axis of an electric dipole})$$

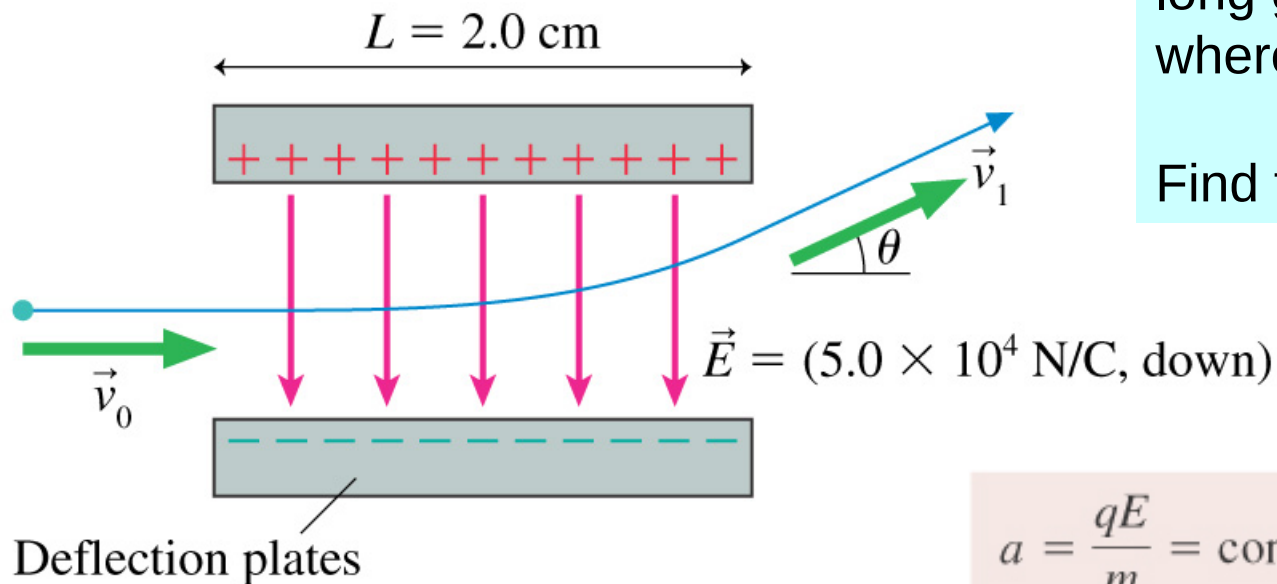
A Dipole

- *Dipole* has zero net charge $\Rightarrow Q=0$, but $\mathbf{p}$ is finite

These dependencies have similar structure but different distance ($\sim r^n$ where $n=1,2,3$) dependencies. Any ideas why is this?

How electric fields acts on charges

Deflecting an Electron Beam



An electron gun creates a beam of electrons moving horizontally: $v_0 = 3.34 \times 10^7 \text{ m/s}$. The electrons enter a 2.0-cm-long gap between electrodes where $E = 5.0 \times 10^4 \text{ N/C}$.

Find the deflection angle θ .

$$a = \frac{qE}{m} = \text{constant} \quad (26.31)$$

Solution

Along x -axis we have motion with constant velocity: $v_x = v_0$

This allows determining time-of-flight: $\Delta t = L/v_x$

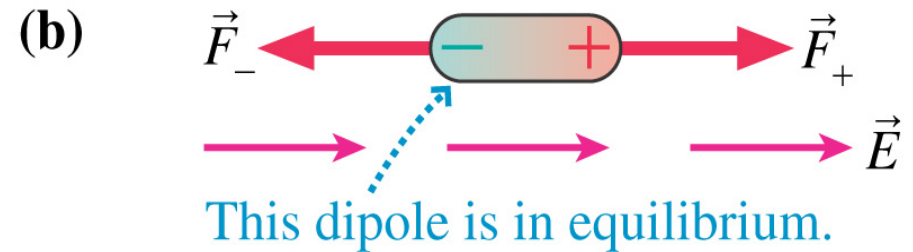
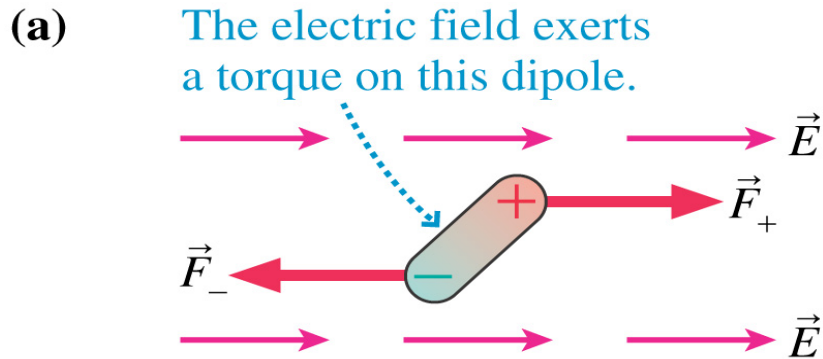
Along y -axis we have motion with constant acceleration:
 $a = F/m$, where $F = eE$

This allows determining vertical component of the resultant velocity:

$$v_y = a \Delta t = eEL/mv_0$$

Finally, $\theta = \tan^{-1} (v_y/v_x)$

How electric field acts on a dipole



- $\mathbf{E}$ is *not* the field of the dipole, but, instead, is a field to which the dipole is responding!

$$\mathbf{F}_+ = +q\mathbf{E} \text{ and } \mathbf{F}_- = -q\mathbf{E}$$

- These forces are equal but opposite, thus:

$$\mathbf{F}_{\text{net}} = \mathbf{F}_+ + \mathbf{F}_- = 0$$

- Because two forces are *not aligned*, the field causes the dipole to *rotate*.

- What is the torque on this dipole

- Two forces are *aligned*.
- What is the torque on this dipole?

To remind about rotational motion

Striking Analogy

Newton's law for straight motion: $\mathbf{a} = \mathbf{F}/m$
 $\mathbf{a}$ - acceleration
 $\mathbf{F}$ - Force
 m - mass

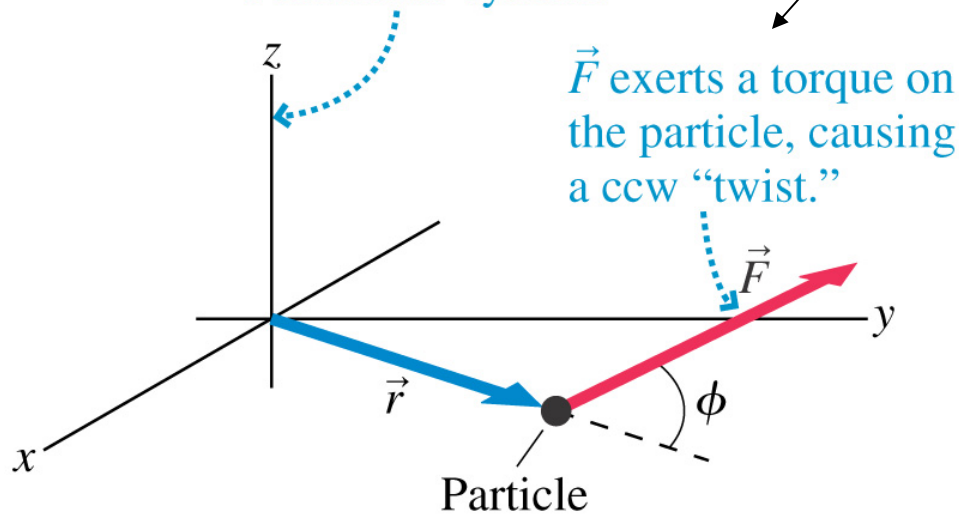


Newton's law for rotational motion: $\alpha = \tau_{\text{net}}/I$
 α - angular acceleration
 $\tau = \mathbf{r} \times \mathbf{F}$ - torque
 $I = \sum m_i r_i^2$ - moment of inertia

Torque (For a review see Chapter 13)

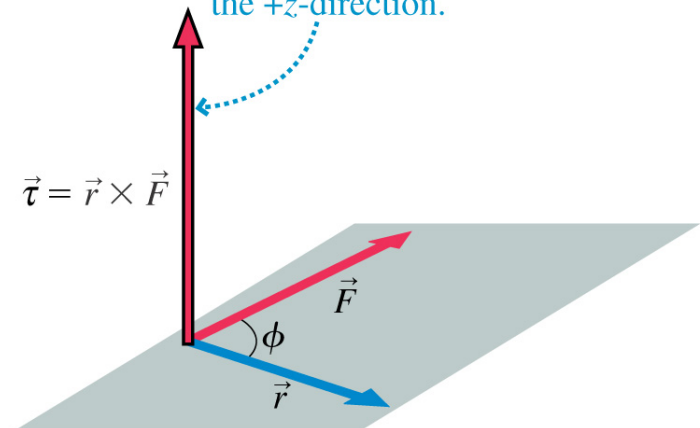
(a)

This is a right-handed coordinate system.



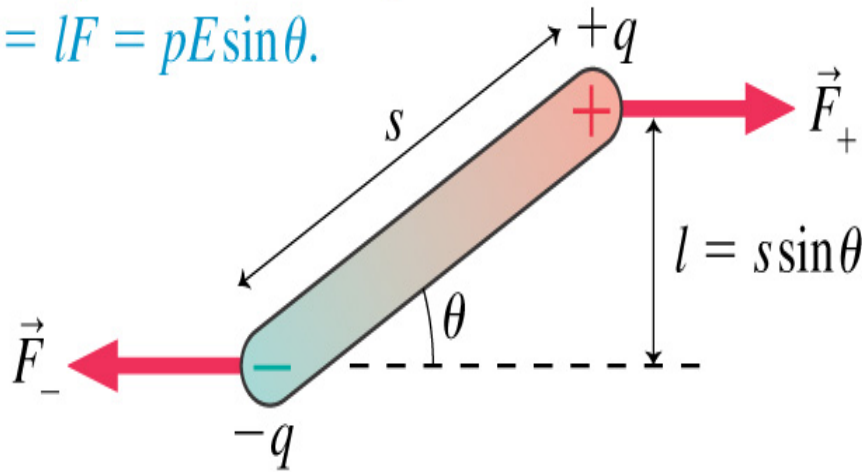
(b)

The torque vector is perpendicular to the plane of $\vec{r}$ and $\vec{F}$, pointing in the $+z$ -direction.



How to calculate the torque on a dipole?

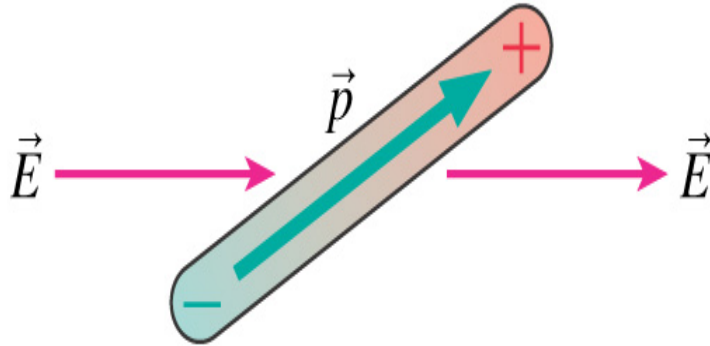
The torque due to a couple is $\tau = lF = pE \sin \theta$.



Torque (τ) is a product of F with the distance (l) between the forces lines (Chapter 13):

$$\tau = Fl = (qE)(s \sin \theta) = pE \sin \theta$$

In terms of vectors, $\vec{\tau} = \vec{p} \times \vec{E}$.



End of Lecture 4
Reading: Entire Chapter 26
Review for Quiz 2
Home Work 1 and 2 in
Mastering Physics